

CHEMISTRY 273

PROJECTION OPERATORS

To determine the symmetry-adapted linear combinations (SALC's) of equivalent a.o.'s from a given basis set, we have been secretly using a simple example of a projection operator. In general, a projection operator is defined by the characteristic property of idempotency:

$$(1) \quad \hat{\sigma}\hat{\sigma}f = \hat{\sigma}f \quad \text{or} \quad \hat{\sigma}^2 = \hat{\sigma}$$

A projection operator may also be thought of as an operator which operates on a space spanned by the basis set (f_1, \dots, f_n) with the following results:

$$(2) \quad \hat{\sigma}^i f_j = f_j \delta_{ij}$$

i.e., the i th projection operator takes f_i into itself but annihilates (takes to zero) any other basis function. Thus if a general function can be expanded in terms of the basis:

$$(3) \quad g = \sum_j c_j f_j$$

we can recover the i th basis function by application of the projection operator $\hat{\sigma}^i$:

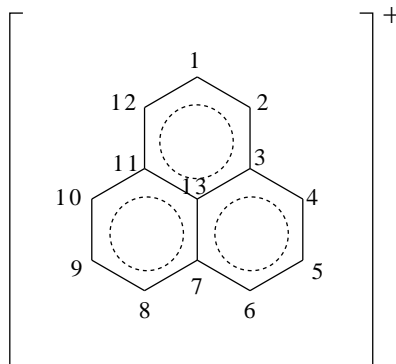
$$(4) \quad \hat{\sigma}^i g = c_i f_i$$

We can use the SALC's of the perinaphthenyl cation as a concrete illustration of projection operators. The three SALC's obtained from the equivalent a.o.'s ϕ_1, ϕ_5 and ϕ_9 are:

$$a'' = (\sqrt{1/3})(\phi_1 + \phi_5 + \phi_9)$$

$$e''_a = (\sqrt{1/6})(2\phi_1 - \phi_5 - \phi_9)$$

$$e''_b = (\sqrt{1/2})(\phi_5 - \phi_9)$$



We will demonstrate that the projection operator for a function in the i th irreducible representation is:

$$(5) \quad \hat{P}^{(i)} = \frac{\ell_i}{h} \sum_R \chi_i^*(R) \hat{R} = \frac{\ell_i}{h} \sum_R \sum_m (\Gamma_i(R))_{mm}^* \hat{R}$$

Letting $\hat{P}^{(i)}$ operate on f_s^j the s th basis function for the j th irreducible representation, we obtain:

$$(6) \quad \begin{aligned} \hat{P}^i f_s^j &= \frac{\ell_i}{h} \sum_R \sum_{m=1}^{\ell_i} (\Gamma_i(R))_{mm}^* \hat{R} f_s^j \\ &= \frac{\ell_i}{h} \sum_R \sum_m (\Gamma_i(R))_{mm}^* \left[\sum_i (\Gamma_j(R))_{st} f_t^j \right] \end{aligned}$$

Since
$$\hat{R} f_s^j = \sum_{t=1}^{\ell_j} (\Gamma_j(R))_{st} f_t^j$$

Thus:

$$(7) \quad \hat{P} f_s^j = \sum_m \sum_t f_t^j \sum_R (\Gamma_i(R))_{mm}^* (\Gamma_j(R))_{st} \frac{\ell_i}{h}$$

by the great orthogonality theorem

$$(8) \quad \hat{P}^i f_s^j = \sum_m \sum_t f_t^j \delta_{ij} \delta_{ms} \delta_{mt}$$

$$(9) \quad \hat{P}^i f_s^j = \sum_m f_m^j \delta_{ij} \delta_{ms} = f_s^j \delta_{ij}$$

Thus $\hat{P}^i f_s^i = f_s^i$ and $\hat{P}^i f_s^j = 0$, $i \neq j$ as required; and \hat{P}^i operating on any function will give us either a function which transforms as the i th irreducible representation or will give us zero (if c_i in the expansion of the general function g is zero).

As an example of this operation, let us look at the application of symmetry projection operators to the a.o. ϕ_1 (here ϕ_1 is the “general function”).

The a.o. can be written as the following linear combination of SALC's:

$$(10) \quad \phi_1 = (\sqrt{3}/3)a_2'' + (\sqrt{6}/3)e_a''$$

The projection $\hat{P}^{a_1''} = (1/12) (\hat{E} + \hat{C}_3 + \hat{C}_3^2 + \hat{C}_2' + \hat{C}_2'' + \hat{C}_2''' - \hat{\sigma}_h - \hat{S}_3 - \hat{S}_3^2 - \hat{\sigma}_v' - \hat{\sigma}_v'' - \hat{\sigma}_v''')$ on ϕ_I gives:

$$\begin{aligned}\hat{P}^{a_1''} \phi_I &= (1/12) (\phi_I + \phi_5 + \phi_9 - \phi_I - \phi_5 - \phi_9 + \phi_I + \phi_5 + \phi_9 - \phi_I - \phi_5 - \phi_9) \\ &= 0 \quad \text{as required!}\end{aligned}$$

The projection $\hat{P}^{a_2''} = (1/12) (\hat{E} + \hat{C}_3 + \hat{C}_3^2 + \hat{C}_2' + \hat{C}_2'' + \hat{C}_2''' - \hat{\sigma}_h - \hat{S}_3 - \hat{S}_3^2 - \hat{\sigma}_v' - \hat{\sigma}_v'' - \hat{\sigma}_v''')$ on ϕ_I gives:

$$\begin{aligned}\hat{P}^{a_2''} \phi_I &= (1/12) (\phi_I + \phi_5 + \phi_9 + \phi_I + \phi_5 + \phi_9 + \phi_I + \phi_5 + \phi_9 + \phi_I + \phi_5 + \phi_9) \\ &= (1/3)(\phi_I + \phi_5 + \phi_9) = (\sqrt{3}/3) a_2'' \quad \text{as required!}\end{aligned}$$

And the projection $\hat{P}^{e''} = (2/12) (2\hat{E} - \hat{C}_3 - \hat{C}_3^2 - 2\hat{\sigma}_h + \hat{S}_3 + \hat{S}_3^2)$ on ϕ_I gives:

$$\begin{aligned}\hat{P}^{e''} \phi_I &= (2/12) (2\phi_I - \phi_5 - \phi_9 + 2\phi_I - \phi_5 - \phi_9) \\ &= (1/3) (2\phi_I - \phi_5 - \phi_9) = (\sqrt{6}/3) e'' \quad \text{as required!}\end{aligned}$$

Thus we demonstrate that SALC's can be generated by applying symmetry projection operators to an a.o.

This procedure is very well defined in two instances: (1) when a given type of a.o. has only one SALC of a given symmetry and (2) when 1 is satisfied and we desire the SALC for a one-dimensional irreducible representation or a single component of a multi-dimensional irreducible representation. If (1) is not satisfied (e.g., the $\phi_2, \phi_4, \phi_6, \phi_8, \phi_{10}, \phi_{12}$ a.o.'s—equivalent in the perinaphthenyl cation have two e'' irreducible representations), the SALC's must be determined by inspection of the actual matrices (not just the characters) for the symmetry operations. If (1) is satisfied and one desires additional SALC's for multidimensional representations, one may proceed in two ways.

The first method utilizes Abelian subgroups of the full symmetry group to generate projection operators. Thus for the D_{3h} group of perinaphthenyl cation we would try the C_{3h} Abelian subgroup.

Using the complex characters of e'' in C_{3h} we get the two projection operators:

$$\hat{P}^{e_+''} = (1/6)(\hat{E} + \varepsilon\hat{C}_3 + \varepsilon * \hat{C}_3^2 - \hat{\sigma}_h - \varepsilon\hat{S} - \varepsilon * \hat{S}_3^2)$$

$$\hat{P}^{e_-''} = (1/6)(\hat{E} + \varepsilon * \hat{C}_3 + \varepsilon\hat{C}_3^2 - \hat{\sigma}_h - \varepsilon * \hat{S} - \varepsilon\hat{S}_3^2)$$

Applying these projection operators to ϕ_1 , one obtains:

$$e_+'' = (1/6)(2\phi_1 + 2\varepsilon\phi_5 + 2\varepsilon^*\phi_9)$$

$$e_-'' = (1/6)(2\phi_1 + 2\varepsilon^*\phi_5 + 2\varepsilon\phi_9)$$

with $\varepsilon = \exp(2\pi i/3)$.

Using $\varepsilon + \varepsilon^* = -1$ and $(\varepsilon - \varepsilon^*)/i = \sqrt{3}$ we can combine e_+'' and e_-'' to get real SALC's for the e'' irreducible representation:

$$e_+'' + e_-'' = (1/3)(2\phi_1 - \phi_5 - \phi_9)$$

$$(1/i)(e_+'' - e_-'') = (\sqrt{3}/3)(\phi_5 - \phi_9)$$

When properly normalized, these are just e_a'' and e_b'' on page 1.

A second method for obtaining components of multidimensional SALC's is to apply the D_{3h} symmetry projection operators to each of the equivalent a.o.'s in a set.

Applying $\hat{P}^{e''}$ to ϕ_5 and ϕ_9 would get:

$$\hat{P}^{e''}\phi_5 = (1/3)(2\phi_5 - \phi_9 - \phi_1)$$

$$\hat{P}^{e''}\phi_9 = (1/3)(2\phi_9 - \phi_1 - \phi_5)$$

These two SALC's plus e_a'' are not linearly independent. However, one can take a linear combination of the two vectors above to get e_b'' as on page 1. The foolproof (take note) method for producing a set of orthogonal (therefore linearly independent) functions from a non-orthogonal set is an application of the Schmidt orthogonalization procedure, which is the subject for another day!