

CHEMISTRY 273
SPECTROSCOPY
(BOHM CH. 18, SEC. 16-20)

$$1. \quad \hat{\mathcal{H}}\Psi(\vec{r}, \vec{\sigma}_e; \vec{R}_N, \vec{\sigma}_N, t) = -\frac{\hbar}{i} \frac{\partial \Psi}{\partial t}$$

(a) If $\hat{\mathcal{H}}$ is independent of t

$$\Psi = \psi(\vec{r}, \vec{\sigma}_e; \vec{R}_N, \vec{\sigma}_N) e^{-iEt/\hbar} \quad \hat{H}\psi = E\psi$$

(b) If $\hat{\mathcal{H}}$ depends on t use time dependent perturbation theory

$$\hat{\mathcal{H}} = \hat{H}^0 + \hat{H}'(t)$$

$$\Psi_\alpha^0 = \psi_\alpha^0 e^{-i(E_\alpha^0/\hbar)t}$$

$$\hat{H}^0 \Psi_\alpha^0 = -\frac{\hbar}{i} \frac{\partial \Psi_\alpha^0}{\partial t} \quad \hat{H}^0 \psi_\alpha^0 = E_\alpha^0 \psi_\alpha^0$$

$$\text{assume expansion } \Psi = \sum_\alpha c_\alpha(t) \Psi_\alpha^0 = \sum_\alpha c_\alpha(t) \psi_\alpha^0 e^{-i(E_\alpha^0/\hbar)t}$$

$$\text{Find } c_\alpha \text{'s by substituting } \hat{\mathcal{H}} \Psi = -\frac{\hbar}{i} \frac{\partial \Psi}{\partial t}$$

$$\hat{H}^0 \left(\sum_\alpha c_\alpha(t) \Psi_\alpha^0 \right) + \hat{H}' \left(\sum_\alpha c_\alpha(t) \Psi_\alpha^0 \right) = -\frac{\hbar}{i} \frac{\partial}{\partial t} \left(\sum_\alpha c_\alpha(t) \Psi_\alpha^0 \right)$$

$$\hat{H}^0 \left(\sum_\alpha c_\alpha(t) \Psi_\alpha^0 \right) + \hat{H}' \left(\sum_\alpha c_\alpha(t) \Psi_\alpha^0 \right) = -\frac{\hbar}{i} \sum_\alpha \left(\frac{\partial c_\alpha}{\partial t} \Psi_\alpha^0 \right) + \sum_\alpha c_\alpha \left(-\frac{\hbar}{i} \frac{\partial \Psi_\alpha^0}{\partial t} \right)$$

I

II

III

IV

I = IV by above

$$\sum_\alpha c_\alpha(t) \hat{H}'(t) \Psi_\alpha^0 = -\frac{\hbar}{i} \sum_\alpha \Psi_\alpha^0 \frac{\partial c_\alpha}{\partial t}$$

II

=

III

multiply by $(\Psi_\beta^0)^*$ and integrate over space

$$\sum_{\alpha} c_{\alpha}(t) \langle \Psi_{\beta}^0 | \hat{H}'(t) | \Psi_{\alpha}^0 \rangle = -\frac{\hbar}{i} \sum_{\alpha} \langle \Psi_{\beta}^0 | \Psi_{\alpha}^0 \rangle \frac{\partial c_{\alpha}}{\partial t}$$

$$\left(\Psi_{\beta}^0 e^{-i(E_{\beta}^0/\hbar)t} \right)^* \quad \Psi_{\alpha}^0 e^{-i(E_{\alpha}^0/\hbar)t} \quad \delta_{\alpha\beta}$$

$$\sum_{\alpha} c_{\alpha}(t) e^{i(E_{\beta}^0 - E_{\alpha}^0)t/\hbar} H'_{\beta\alpha} = -\frac{\hbar}{i} \frac{\partial c_{\beta}}{\partial t}$$

$$\int_{\text{space}} (\Psi_{\beta}^0)^* \hat{H}'(t) \Psi_{\alpha}^0 d\tau$$

system starts in state α at time $t = t_0$

then $c_{\alpha}(t_0) = 1$ $c_{\alpha \neq \beta}(t_0) = 0$

at $t = t_0$ $-\frac{\hbar}{i} \frac{\partial c_{\beta}}{\partial t} = e^{i(E_{\beta}^0 - E_{\alpha}^0)t/\hbar} H'_{\beta\alpha}$

at later time $c_{\beta} = -\frac{i}{\hbar} \int_{t_0}^t e^{i\frac{(E_{\beta}^0 - E_{\alpha}^0)}{\hbar}t''} H'_{\beta\alpha} dt''$

I. In weak electromagnetic field

$$\hat{H}' = -\frac{e}{2mc} (\vec{A} \cdot \vec{p} + \vec{p} \cdot \vec{A}) \text{ where } \vec{A} \text{ is vector potential:}$$

wave in z-direction with electric field in x direction.

$$\vec{A} = A_0 \cos(kz - \omega t) \vec{i} \qquad \nabla \times \vec{A} = \vec{B} = \mu \vec{H}$$

$$\omega = 2\pi\nu \qquad \vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$k = \frac{2\pi}{\lambda}$$

or kept explicitly in terms of the electric and magnetic field interactions

$$\psi^* \vec{\mathcal{E}} \psi \text{ and } \psi^* (\vec{v} \times \vec{B}) \psi$$

see McQ. p. 472-475.

$$\vec{B} = -A_o k \sin(kz - \omega t) \vec{j} = -A_o \frac{\omega}{c} \sin(kz - \omega t) \vec{j}$$

$$\vec{\mathcal{E}} = -A_o \frac{\omega}{c} \sin(kz - \omega t) \vec{i}$$

Result #1 for single frequency ν

$$|c_\beta|^2 = \frac{e^2}{m^2 c^2} A_o^2 |Q_{\alpha\beta}|^2 \sin^2 \left[\frac{(E_\beta - E_\alpha - h\nu)(t - t_o)}{2k} \right] (E_\beta - E_\alpha - h\nu)^2$$

$$Q_{\alpha\beta} = \int \psi_\alpha^o \hat{p}_x e^{ikz} \psi_\beta^o d\tau$$

Result #2 for a small spread of frequencies with intensities $I(\nu)$ centered around $\nu_o \approx \nu_{\alpha\beta}$

$$|c_\beta|^2 = \frac{2\pi e^2}{m^2 c \hbar^2} \frac{I(\nu_o)}{\omega_{\beta\alpha}} |Q_{\beta\alpha}|^2 (t - t_o)$$

[see handout for proofs of above arguments]

Probability of transition per unit time

$$P_{\beta \leftarrow \alpha} = \frac{2\pi e^2 |Q_{\beta\alpha}|^2}{\omega_{\alpha\beta}^2 m^2 c \hbar^2} \quad \left(\begin{array}{l} \text{probability for x-polarized} \\ \text{light traveling in z-direction.} \end{array} \right)$$

$$Q_{\beta\alpha} = \int (\psi_\beta^o)^* \left(-\frac{\hbar}{i} \frac{\partial}{\partial x} e^{ikz} \psi_\alpha^o \right) d\tau$$

$$Q_{\beta\alpha} = \int (e^{ikz}) (\psi_\beta^o)^* \hat{p}_x \psi_\alpha^o d\tau$$

for molecule at $z = z_0$

$$Q_{\beta\alpha} = \int \left[1 + ik(z - z_0) - \frac{k^2}{2}(z - z_0)^2 \dots \right] (\psi_\beta^0)^* \hat{p}_x \psi_\alpha^0 d\tau$$

- 1) Dipole approximation: extent of molecule so small compared to wavelength of light, $k(z - z_0)$ is small in region of interest. $\left(k = \frac{2\pi}{\lambda} \right)$

$$Q_{\beta\alpha} \approx \int \psi_\beta^0 \hat{p}_x \psi_\alpha^0 d\tau \quad \text{momentum transition moment}$$

$$\int \psi_\beta^0 \hat{p}_x \psi_\alpha^0 d\tau = i\omega_{\beta\alpha} m \int \psi_\beta^0 \hat{x} \psi_\alpha^0 d\tau \quad \text{position transition moment [see handout]}$$

$$= \frac{-i\omega_{\beta\alpha}}{e} m \mu_{\beta\alpha} \quad \langle \vec{\mu} \rangle_{\beta\alpha} = \int \psi_\beta (-e\vec{r}) \psi_\alpha d\tau$$

$$P_{\beta\alpha} = \frac{2\pi}{c\hbar^2} \langle \mu_x \rangle^2$$

for random alignment of light and molecule

$$\bar{P}_{\beta\alpha} = \frac{2}{3} \frac{\pi}{c\hbar^2} \left[\langle \mu_x \rangle^2 + \langle \mu_y \rangle^2 + \langle \mu_z \rangle^2 \right]$$

$$f = \frac{2\nu_0 \hbar m c}{e^2} \bar{P}_{\alpha\beta} = 4.315 \times 10^{-9} \int \epsilon \, d\nu$$

||

mol⁻¹ L cm⁻¹ cm⁻¹

oscillator strength (unitless) $\frac{4\pi\nu_0 m}{3\hbar e^2} |\vec{\mu}_{\beta\alpha}|^2$

$$P_{\beta\alpha} = \frac{2}{3} \frac{\pi}{c\hbar^2} |\vec{\mu}_{\beta\alpha}|^2 \quad \vec{\mu}_{\beta\alpha} = \int \psi_\beta^0 (-e\vec{r}) \psi_\alpha^0 d\tau$$

for one electron.